

Calcular los símbolos de Christoffel y verificar el valor de R .

Tenemos el intervalo ds^2 por, usamos notación en t , para:

$$ds^2 = -(1+2\phi) dt^2 + a^2(1-2\psi) \delta_{ij} dx^i dx^j$$

Lo que nos da las componentes de nuestro tensor métrico

$$g_{00} = a^2(2\psi+1) \quad g^{00} = \frac{1}{a^2} \frac{1}{2\psi+1} \sim \frac{1}{a^2} (1-2\psi)$$

$$g_{ij} = a^2(2\phi-1) \delta_{ij} \quad g^{ij} = \frac{1}{a^2} \frac{1}{2\phi-1} \sim \frac{-1}{a^2} (1+2\phi) \delta^{ij}$$

$$\Gamma_{\mu\nu}^\alpha = \frac{1}{2} g^{\alpha\rho} (g_{\rho\mu,\nu} + g_{\rho\nu,\mu} - g_{\mu\nu,\rho})$$

$$\Gamma_{00}^0 = \frac{1}{2} g^{0\rho} (g_{\rho 0,0} + g_{\rho 0,0} - g_{00,\rho}) \quad , \text{ debe } \rho=0$$

$$= \frac{1}{2} g^{00} (g_{00,0} + g_{00,0}) = \frac{1}{2} g^{00} 2 (\partial_0 a^2(2\psi+1))$$

$$= g^{00} (2a\dot{a}(2\psi+1) + a^2(2\dot{\psi}+1))$$

$$= \frac{(1-2\psi)}{2a^2} (2a\dot{a}(2\psi+1) + a^2(2\dot{\psi}+1))$$

$$= (1-2\psi) \frac{\dot{a}}{a} (2\psi+1) + \frac{1}{2} (2\dot{\psi}+1) = (1-2\psi) \left[2(2\psi+1) + \frac{1}{2}(2\dot{\psi}+1) \right]$$

Eliminamos los términos con ψ y sólo queda el de \mathcal{H} y $\mathcal{H}\psi$

$$= \cancel{2(1-2\psi)} \cancel{2(1-2\psi)} + \frac{1}{2} 2\dot{\psi} = \boxed{2(1+\dot{\psi})}$$

$$\Gamma_{0i}^0 = \frac{1}{2} g^{00} (g_{00,i} + \cancel{g_{i0,0}} - \cancel{g_{0i,0}})$$

$$= \frac{1}{2} g^{00} (\partial_i [a^2 (2\psi + 1)]) , \partial_i a^2 = 0$$

$$= \frac{1}{2} g^{00} 2(a^2 \partial_i \psi) = g^{00} \partial_i \psi a^2 = \frac{1}{a^2} (1 - 2\psi) \partial_i \psi a^2$$

$$= \partial_i \psi - 2\psi \cancel{\partial_i \psi} \rightarrow 0(2)$$

$$= \boxed{\partial_i \psi}$$

$$\Gamma_{00}^i = \frac{1}{2} g^{i\beta} (g_{\beta 0,0} + g_{\beta 0,0} - g_{00,\beta}) , \text{ pero } i=\beta$$

$$= \frac{1}{2} g^{ii} (-g_{00,i})$$

$$= \frac{1}{2} (-\frac{1}{a^2} (2\psi + 1)) (-\partial_i (a^2 (2\psi + 1))) , \partial_i a^2 = 0$$

$$= \frac{1}{2} (2\psi + 1) \partial_i (2\psi + 1)$$

$$= \frac{1}{2} (2\psi + 1) (2\partial_i \psi) \quad \sim \text{los términos cruzados los eliminamos}$$

$$= \frac{1}{2} 2 \partial_i \psi = \underline{\partial_i \psi}$$

$$= \boxed{\delta^{ij} \partial_j \psi}$$

Ahora

$$\Gamma_{ij}^0 = \frac{1}{2} g^{0\beta} (g_{\beta i,j} + g_{\beta j,i} - g_{ij,\beta}) , \beta=0$$

$$= \frac{1}{2} g^{00} (g_{ij,0})$$

$$= \frac{1}{2} g^{00} [\partial_t (a^2 (2\psi - 1) \delta_{ij})]$$

$$\begin{aligned}
&= \frac{1}{2} g^{00} (2a\dot{a} (2\phi - 1) + 2a^2 \dot{\phi}) \delta_{ij} \\
&= \frac{1}{2} \frac{1}{a^2} (1 - 2\psi) (2a\dot{a} (2\phi - 1) + 2a^2 \dot{\phi}) \delta_{ij} \\
&= (1 - 2\psi) (2(2\phi - 1) + \dot{\phi}) \delta_{ij} \\
&= (2H + 22H\phi - 2\psi H - \dot{\phi}) \delta_{ij} \\
&= \boxed{2H\delta_{ij} - [\dot{\phi} + 2H(\psi + \phi)] \delta_{ij}}
\end{aligned}$$

Ahora $\Gamma_{j0}^i = \frac{1}{2} g^{i\beta} (g_{\beta j,0} + g_{\beta 0,j} - g_{j0,\beta}) \quad \beta = i$

$$\begin{aligned}
&= \frac{1}{2} g^{ii} (g_{ij,0} + \cancel{g_{i0,j}} - \cancel{g_{j0,i}}) \\
&= \frac{1}{2} g^{ii} (g_{ij,0}) = \frac{1}{2} - \frac{(1+2\phi)}{a^2} (2a\dot{a}(2\phi-1)\delta_{ij}) \\
&= -\frac{(1+\phi)}{2a^2} (2a\dot{a}(2\phi-1+2a^2\dot{\phi})\delta_{ij}) \\
&= -(1+\phi) (2(2\phi-1) + \dot{\phi}) \delta_{ij} \quad , \text{ solo términos } \mathcal{O}(1) \\
&= \boxed{(H - \dot{\phi}) \delta_{ij}}
\end{aligned}$$

Por último

$$\begin{aligned}
\Gamma_{ij}^{\kappa} &= \frac{1}{2} g^{\kappa\beta} (g_{\beta i,j} + g_{\beta j,i} + g_{ij,\beta}) \quad , \quad \kappa = \beta \\
&= \frac{1}{2} g^{\kappa\kappa} (g_{\kappa i,i} + g_{\kappa j,i} + g_{ij,\kappa})
\end{aligned}$$

$$= -\frac{(1+2\phi)}{a^2} \left[\delta_{ij} a^2 (\partial_\kappa (2\phi - 1)) + \delta_{i\kappa} a^2 \partial_j (2\phi - 1) - \delta_{j\kappa} a^2 \partial_i (2\phi - 1) \right]$$

$$= \cancel{1+2\phi} (\delta_{ij} \phi_{,\kappa} - \delta_{i\kappa} \phi_{,j} + \delta_{j\kappa} \phi_{,i})$$

$$= -\delta_{ij} \phi_{,\kappa} - \delta_{i\kappa} \phi_{,j} + \delta_{j\kappa} \phi_{,i}$$

agregando una f^{il} y agrupando queda igual que en la tarea

Para calcular R , tenemos

$$R = g^{\mu\nu} R_{\mu\nu}$$

y podemos escribirlo como parte temp + parte esp

$$R = g^{00} R_{00} + g^{ij} R_{ij}$$

Ya tenemos g^{00} y g^{ij} , vamos a calcular R_{00}

$$\begin{aligned} \Rightarrow R_{00} &= \Gamma_{00,\mu}^{\mu} - \Gamma_{0\mu,0}^{\mu} + \Gamma_{\ell\mu}^{\mu} \Gamma_{00}^{\ell} - \Gamma_{\ell 0}^{\mu} \Gamma_{0\mu}^{\ell} \\ &= \Gamma_{00,i}^i - \Gamma_{0i,0}^i + \Gamma_{0i}^i \Gamma_{00}^0 - \cancel{\Gamma_{00}^i \Gamma_{0i}^0} + \cancel{\Gamma_{0i}^i \Gamma_{00}^0} - \Gamma_{i0}^i \Gamma_{0i}^0 \end{aligned}$$

ambos contienen perturbación $\mathcal{O}(1)$

$$= f^{ij} \lambda_i \partial_j \psi - \partial_t (\epsilon (\delta_j^i - \phi \delta_j^i)) + \partial_i \psi (\epsilon (1 - \phi)) - (4 - \phi) \delta_i^j \delta_j^i$$

$$\begin{aligned}
 &= \cancel{\partial_{ii}} \nabla^2 \Psi - 3 \partial_t (\mathcal{H} - \dot{\phi}) - 3(\mathcal{H} + \dot{\psi})(\mathcal{H} - \dot{\phi}) - (\mathcal{H} - \dot{\phi}) \delta_{ij} \delta_{ij} \\
 &= \nabla^2 \Psi - 3\dot{\mathcal{H}} + 3\ddot{\phi} - \cancel{3\mathcal{H}^2 + 3\mathcal{H}^2} + 3\mathcal{H}(\dot{\phi} + \dot{\psi}) \\
 &= \nabla^2 \Psi - 3\dot{\mathcal{H}} + 3\ddot{\phi} + 3\mathcal{H}(\dot{\phi} + \dot{\psi})
 \end{aligned}$$

Now $R_{ij} = \Gamma_{ij,\mu}^{\mu} - \Gamma_{i\mu,\mu}^{\mu} + \Gamma_{\ell\mu}^{\mu} \Gamma_{ij}^{\ell} - \Gamma_{\ell j}^{\mu} \Gamma_{i\mu}^{\ell}$

$$\begin{aligned}
 &= \Gamma_{ij,0}^0 + \Gamma_{ij,\kappa}^{\kappa} - \Gamma_{i0,j}^0 - \Gamma_{\kappa\kappa,i}^{\kappa} + \Gamma_{00}^0 \Gamma_{ij}^0 + \Gamma_{0\kappa}^{\kappa} \Gamma_{ij}^0 + \Gamma_{\kappa j}^0 \Gamma_{i0}^{\kappa} \\
 &\quad - \Gamma_{0j}^{\kappa} \Gamma_{i\kappa}^0 \\
 &= \delta_{ij} a^2 \left(\mathcal{H}^2 - 2\mathcal{H}^2(\phi + \psi) + 4\mathcal{H}\dot{\phi} + \frac{\ddot{a}}{a} - 2\frac{\ddot{a}}{a}(\phi + \psi) - \mathcal{H}(\dot{\psi}) - \ddot{\phi} \right) \\
 &\quad - 2\partial_i \partial_j \phi + \delta_{ij} \partial_{\kappa} \partial_{\kappa} \phi - \partial_i \partial_j \psi + \delta_{ij} \partial_{\kappa} \partial_{\kappa} \phi \\
 &\quad + \delta_{ij} a^2 \mathcal{H} \ddot{\phi} + \delta_{ij} a^2 [\mathcal{H}^2 - 2\mathcal{H}^2(\phi + \psi) - 2\mathcal{H}\dot{\phi}]
 \end{aligned}$$

Agrupamos con las δ y las parciales

$$\begin{aligned}
 &= \delta_{ij} a^2 \left\{ \frac{1}{a^2} \nabla^2 \phi + \frac{\ddot{a}}{a} [1 - 2(\psi + \phi)] - \mathcal{H}\dot{\psi} - \ddot{\phi} + 2\mathcal{H}^2 \right. \\
 &\quad \left. - 4\mathcal{H}^2(\psi + \phi) - 6\mathcal{H}\dot{\phi} \right\} + (\phi - \psi) \delta_{ij} \\
 &= \delta_{ij} a^2 \left\{ \frac{1}{a^2} \nabla^2 \phi + \mathcal{H} [1 - 2(\psi + \phi)] - \mathcal{H}(\dot{\psi} - \dot{\phi}) + 2\mathcal{H}^2 \right. \\
 &\quad \left. + 2\mathcal{H}^2 - 4\mathcal{H}^2(\psi + \phi) + 6\mathcal{H}\dot{\phi} \right\} + (\phi - \psi) \delta_{ij}
 \end{aligned}$$

$$\begin{aligned}
 R_{ij} = & \delta_{ij} \left\{ \nabla^2 \phi - \ddot{\phi} + [1 - 2(\phi + \psi)](\mathcal{H} + 2\mathcal{H}^2) - \mathcal{H}(\dot{\phi} + 5\dot{\phi}) \right\} \\
 & - (\phi - \psi) \delta_{ij}
 \end{aligned}$$

⇒ Allora potremmo avere

$$R = g^{00} R_{00} + g^{ij} R_{ij}$$

$$= \underline{(1-2\psi) R_{00}} - (1+2\phi) \delta^{ij} R_{ij}$$

$$= (1-2\psi) [\nabla^2 \psi - 3\dot{H} + 3\ddot{\phi} + 3H(\dot{\phi} + \dot{\psi})]$$

$$= \nabla^2 \psi - 3\dot{H} + 3\ddot{\phi} + 3H(\dot{\phi} + \dot{\psi}) - 6\dot{H}\psi$$

$$= \underline{\nabla^2 \psi - 3\dot{H}(1+2\dot{H}\psi) + 3\ddot{\phi} + 3H(\dot{\phi} + \dot{\psi})}$$

Allora $-(1+2\phi) \delta^{ij} R_{ij}$

$$= -(1+2\phi) \delta^{ij} \left[\delta_{ij} \left\{ \nabla^2 \phi - \ddot{\phi} + [1-2(\phi+\psi)](\dot{H}+2\dot{H}^2) - H(\dot{\phi}+5\dot{\phi}) \right\} - (\phi-\psi)_{,ij} \right]$$

$$= -(1+2\phi) \left[3 \left\{ \nabla^2 \phi - \ddot{\phi} + [1-2(\phi+\psi)](\dot{H}+2\dot{H}^2) - H(\dot{\phi}+5\dot{\phi}) \right\} - (\phi-\psi)_{,ii} \right]$$

$$= 3 \left\{ \nabla^2 \phi - \ddot{\phi} + [1-2(\phi+\psi)](\dot{H}+2\dot{H}^2) - H(\dot{\phi}+5\dot{\phi}) \right\} - \nabla^2 (\phi-\psi) + 2\phi \cdot 3 \{ (\dot{H}+2\dot{H}^2) \}$$

$$R = \nabla^2 \psi - 3\dot{H}(1+2\dot{H}\psi) + 3\ddot{\phi} + 3H(\dot{\phi} + \dot{\psi}) + 3 \nabla^2 \phi - \ddot{\phi} + [1-2(\phi+\psi)](\dot{H}+2\dot{H}^2) - H(\dot{\phi}+5\dot{\phi}) - \nabla^2 (\phi-\psi) + 2\phi \cdot 3 \{ (\dot{H}+2\dot{H}^2) \}$$