Updated Cosmology

with Python



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Homework 5.b

1) The most general spherically symmetric metric can be written as

$$ds^{2} = -e^{2F(r,t)}dt^{2} + e^{2H(r,t)}dr^{2} + r^{2}(d\theta^{2} + \sin^{2}\theta d\phi^{2}).$$
(1)

This metric is very important as it underlies the theory of both homogeneous cosmological models and spherically symmetric models for massive stars and black holes. The functions F(r,t) and H(r,t) are determined by the material content of the space-time as described by the energy-momentum tensor and by the boundary conditions defining the problem. Compute the components of the Einstein tensor.

$$G_{00} = e^{-2H} \left(\frac{2}{r}H' - \frac{1}{r^2}\right) + \frac{1}{r^2},$$

$$G_{11} = e^{-2H} \left(\frac{2}{r}F' + \frac{1}{r^2}\right) - \frac{1}{r^2},$$

$$G_{22} = G_{33} = e^{-2H} \left(F'' + F'^2 - H'F' + \frac{1}{r}(F' - H')\right)$$

$$- e^{-2F} \left(\ddot{H} + \dot{H}^2 - \dot{H}\dot{F}\right),$$

$$G_{01} = G_{10} = \frac{2}{r}\dot{H} e^{-(F+H)}.$$

2) The Bianchi models are a large family of homogeneous but anisotropic cosmological models. We consider the homogenous and anisotropic space-time described by Bianchi type-III metric in the form

$$ds^{2} = dt^{2} - A(t)^{2} dx^{2} - B(t)^{2} e^{-2\alpha x} dy^{2} - C(t)^{2} dz^{2},$$
(2)

where A(t), B(t) and C(t) are the scale factors (metric tensors) and functions of the cosmic time t, and α is a constant. (Bianchi type-I metric can be recovered by choosing $\alpha = 0$). Here, we assume an anisotropic fluid whose energy-momentum tensor is in diagonal form:

$$T_{\nu}^{\ \mu} = \text{diag}[1, -w_x, -w_y, -w_z] = \text{diag}[1, -w, -(w+\gamma), -(w+\delta)]\rho, \tag{3}$$

where ρ is the energy density of the fluid, w_x , w_y and w_z are the directional EoS parameters on the x, y and z axes respectively; w is the deviation-free EoS parameter of the fluid. δ and γ are not necessarily constants and can be functions of the cosmic time t.

a) Compute the components of the Einstein's field equations

$$\frac{\dot{A}}{A}\frac{\dot{B}}{B} + \frac{\dot{A}}{A}\frac{\dot{C}}{C} + \frac{\dot{B}}{B}\frac{\dot{C}}{C} - \frac{\alpha^2}{A^2} = \rho,$$
(5)

$$\frac{\ddot{B}}{B} + \frac{\ddot{C}}{C} + \frac{\dot{B}}{B}\frac{\dot{C}}{C} = -w\rho, \qquad (6)$$

$$\frac{\ddot{A}}{A} + \frac{\ddot{C}}{C} + \frac{\dot{A}}{A}\frac{\dot{C}}{C} = -(w+\delta)\rho, \qquad (7)$$

$$\frac{\ddot{A}}{A} + \frac{\ddot{B}}{B} + \frac{\dot{A}}{A}\frac{\dot{B}}{B} - \frac{\alpha^2}{A^2} = -(w+\gamma)\rho,$$
(8)

$$\alpha \left(\frac{\dot{A}}{A} - \frac{\dot{B}}{B}\right) = 0 \tag{9}$$

where the overdot denotes derivation with respect to the cosmic time t.

b) Show the solution of Eqn. (9) gives $B = c_1 A$, where c_1 is the positive constant of integration. c) Substitute this solution into (7), and subtract the result from (6), to show that the skewness parameter on the y axis is null, i.e. $\delta = 0$, which means that the directional EoS parameters, hence the pressures, on the x and y axes are equal.

The directional Hubble parameters in the directions of x, y and z for the Bianchi type-III metric may be defined as follows,

$$H_x \equiv \frac{\dot{A}}{A}, \quad H_y \equiv \frac{\dot{B}}{b}, \quad H_z \equiv \frac{\dot{C}}{C}.$$
 (4)

and the mean Hubble parameter is given by

$$H = \frac{1}{3}\frac{\dot{V}}{V} = \frac{1}{3}\left(\frac{\dot{A}}{A} + \frac{\dot{B}}{B} + \frac{\dot{C}}{C}\right)$$
(5)

By solving the set of field equations, the difference between the expansion rates on x and z axes could be found $H_x - H_z$, [see Gen Relativ Gravit (2010) 42:763–775].

3) The action of a massive scalar field, with potential $V(\phi)$, is given by

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right].$$
 (6)

a) Show that the corresponding field equation for ϕ , obtained from the Euler-Lagrange equations, reads as

$$\Box^2 \phi + \frac{dV}{d\phi} = 0. \tag{7}$$

where the d'Alembertian is : $\Box^2 = \partial_a \partial^a$.