



Tonatiuh Matos http://www.fis.cinvestav.mx/~tmatos/

http://www.iac.edu.mx

The LCDM Some Problems of the LCDM The SFDM Alternative



The Standard Model of Cosmology: LCDM



CDM Paradigme

- 1.- The DM fluctuation spectrum at recombination is detemined by a small number of physical parametters.
- 2.- After recombination, the amplitude of the baryonic fluctuations rapidly grows to match that of the DM fluctuations.
- 3.- Smaller-mass fluctutions grow to nonlineartity and virialize and then are hirarchically clustered within successively larger bound systems.
- 4.- Ordinary matter in bound systems of total mass $10^{8-12} M_{\odot}$ cools rapidly enough within the DM halos to form galaxies, while larger mass fluctutations form clusters



The Standard Model of Cosmology: LCDM







The Standard Model of Cosmology: LCDM







The Standard Model of Cosmology: Problems



- Extreme fine tuning
- Coincidence
- Cuspy central density profiles
- Missing Satellites Problem
- No-fit of the early assembly of galaxies
- Galaxies seems to be simpler
- Voids are too empty
- No-detection of DM
- etc.



The Standard Model of Cosmology: Problems



- Cuspy central density profiles
- Missing Satellites Problem





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Galaxy's Center: Observations



W. J. G. de Blok, Stacy S. McGaugh, Albert Bosma, and Vera C. Rubin. ApJ Letters 552(2001)L24 Cinvestav LOW SURFACE BRIGHTNESS MASS DENSITY PROFILES Vol. 552



 $\log(\rho/[M_{\odot} \text{ pc}^{-3}])$

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Galaxy's Center: Observations





Galaxy's Center: Observations





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F. Gobernato, et. al., NATURE, Vol 463, 14 January 2010







Stellar Mass Predictions



Till Sawala, Qi Guo, Cecilia Scannapieco, Adrian Jenkins and Simon White. Mon. Not. Roy. Astron. Soc. 413, (2011), 659





Stellar Mass Predictions



V. Avila-Reese, P. Colín, A. González-Samaniego, O. Valenzuela, C. Firmani, H. Velázquez, & D. Ceverino. ApJ, 736:134, (2011)



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Velocity predictions





Jesús Zavala, et. al. ApJ, 700:1779-1793, 2009 August 1

The simulation with CDM predicts a steep rise in the VF toward lower velocities; for Vmax = 35 km/s, it forecasts ~10 times more sources than the ones observed. If confirmed by the complete ALFALFA survey, these results indicate a potential problem for the CDM paradigm





M. Vogelsberger, S. Genel, V. Springel, P. Torrey, D. Sijacki, D. Xu, G. Snyder, S. Bird, D. Nelson & L. Hernquist. Nature 509 (2014) 177



Metallicity



Andrew Pontzen & Fabio Governato. Nature 506 (2014) 171





Galaxy Satellites







Metallicity







Metallicity



Evan N. Kirby, Judith G. Cohen, Puragra Guhathakurta, Lucy Cheng, James S. Bullock, Anna Gallazzi. ApJ (2013) 779 102





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Michael Boylan-Kolchin, James S. Bullock and Manoj Kaplinghat. Mon.Not.Roy.Astron.Soc.422:(2012)1203





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Michael Boylan-Kolchin, James S. Bullock and Manoj Kaplinghat. Mon.Not.Roy.Astron.Soc.422:(2012)1203





Galactic Dynamics



Jorge Peñarrubia, A. J. Benson, Matthew G. Walker, Gerard Gilmore, Alan W. McConnachie and Lucio Mayer. MNRAS 406, 1290–1305 (2010)





Voids are too empty

P. J. E. Peebles & A. Nusser, Nature 465(2010)565





It is not clear why objects that might have been assembled in such very different ways, from different ancestral objects, should have had evolutionary tracks that converged to show small dispersions around simple power law forms for their size–luminosity relations.

Nair, P. B., van den Bergh, S. & Abraham, R. G. The environmental dependence of the luminosity-size relation for galaxies. Astrophys. J. 715, 606–622 (2010).



Voids are too empty

P. J. E. Peebles & A. Nusser, Nature 465(2010)565





In short, the general insensitivity of galaxies to their environments is not expected in standard ideas. It would help if galaxies were more rapidly assembled, as they could then evolve as more nearly isolated island universes.



Early Galaxies Compactness



Pieter G. van Dokkum et. al. NATURE, Vol 460, 6 August 2009. The Astrophysical Journal, 677:L5–L8, 2008 April 10





Early Galaxies Compactness



Pieter G. van Dokkum et. al. NATURE, Vol 460, 6 August 2009. The Astrophysical Journal, 677:L5–L8, 2008 April 10





Some Alternatives



SM Extentions

- Self-Interacting DM
- Warm DM
- Super Heavy DM
- Self-Annihilating DM
- Decaying DM
- Extra Symmetries,
- Repulsive DM
- Fuzzy DM
- k-essence
- Scalar Field DM (BEC DM), etc.

GR Extentions

- Scalar Field DM (BEC DM)
- MOND -> MOND+
- f(R)
- Extra Dimensions
- etc.



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A brief Review of the Scalar Field Dark Matter model. Juan Magana, Tonatiuh Matos, Victor Robles, Abril Suarez. arXiv:1201.6107

Here
$$\mathcal{L} = \sqrt{-g} \left[R - \frac{1}{2} (\nabla \Phi)^2 - V(\Phi) \right] - e^{-2\alpha \Phi} F^2 \right]$$

$$V(\Phi) = W_0 + \frac{11}{2!} N^{2'} \Phi^{2} + \frac{1}{2!!} V^{4'} \Phi^{4} \Phi^{4} + \frac{1}{3!} V^{''} \Phi^{3} + \frac{1}{4!} V^{iv} \Phi^{4} + \cdots$$



Bose-Einstein Condensates



Abril Suarez, Victor H. Robles, Tonatiuh Matos . A Review on the Scalar Field/ Bose-Einstein Condensate Dark Matter Model Astrophysics and Space Science Proceedings 38, Chapter 9 (2013) arXiv:1302.0903

$$\Box \Phi + \frac{dV}{d\Phi} = 0$$

$$\ddot{\Phi} + 3H\dot{\Phi} + \frac{dV}{d\Phi} = 0$$

$$k_B T_c = \frac{2\pi\hbar^2}{m^{\frac{5}{3}}} \left(\frac{\rho}{g_{\frac{3}{2}}(1)}\right)^{\frac{2}{3}}$$



The Cosmology









$$\delta\ddot{\Phi} + 3H\delta\dot{\Phi} - \frac{1}{a^2}\nabla^2\delta\Phi + V_{,\Phi\Phi}\,\delta\Phi + 2V_{,\Phi}\,\phi = 0$$

$$\delta \Phi = \sqrt{\hat{\rho}} \cos\left(\frac{mc^2 t}{\hbar} + S\right) \qquad \qquad \vec{v} \equiv \frac{\hbar}{m} \nabla S$$





Abril Suarez and TM MNRAS 311, (2011), 87



$$\begin{aligned} \frac{\partial \hat{\rho}}{\partial t} &+ \frac{1}{a^2} \nabla \cdot (\hat{\rho} \vec{v}) + 3H\hat{\rho} = 0 \\ &- \frac{\hbar}{m} \hat{\rho} \left(\ddot{S} + 3H\dot{S} \right) - \frac{\hbar}{m} \hat{\rho} \dot{S} = 0 \\ \frac{\partial \vec{v}}{\partial t} &+ \frac{1}{a^2} \vec{v} \nabla \cdot \vec{v} + \nabla \phi + \frac{\hbar^2}{2m^2} \nabla \left(\frac{\Box \sqrt{\hat{\rho}}}{\sqrt{\hat{\rho}}} \right) = 0, \\ &- \frac{\hbar}{m} \dot{S} \frac{\partial \vec{v}}{\partial t} = 0, \end{aligned}$$

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t

$$\hat{\rho} = \hat{\rho}(t) \qquad S = S(t) \qquad \vec{v} \equiv \frac{n}{m} \nabla S$$

$$\frac{\partial \hat{\rho}}{\partial t} + BH \hat{\rho} \frac{1}{a^2} \nabla 0 \cdot (\hat{\rho} \vec{v}) + 3H \hat{\rho} = 0$$









$$\vec{v} \equiv \frac{\hbar}{m} \nabla S$$

$$\frac{\partial \hat{\rho}}{\partial t} + \frac{1}{a^2} \nabla \cdot (\hat{\rho} \vec{v}) \neq \mathbf{B} H \hat{\rho} = 0$$

$$\frac{\partial \vec{v}}{\partial t} + \frac{1}{a^2} \vec{v} \nabla \cdot \vec{v} + \nabla \phi + \frac{\hbar^2}{2m^2} \nabla \left(\frac{\Box \sqrt{\hat{\rho}}}{\sqrt{\hat{\rho}}} \right) = 0,$$





$$\begin{split} \hat{\rho} &= \hat{\rho}_{0} + \rho_{1}(t) \exp(i\vec{k}\cdot\vec{x}/a) & \vec{v}_{1} = \lambda\vec{k} + \vec{v}_{2} \\ \vec{v} &= \vec{v}_{0} + \vec{v}_{1}(t) \exp(i\vec{k}\cdot\vec{x}/a) & \frac{\partial\rho_{1}}{\partial t} + 3H\rho_{1} + i\frac{\hat{\rho}_{0}}{a^{2}}k^{2}\lambda &= 0 \\ & \frac{\partial\lambda}{\partial t} + i(\frac{v_{q}^{2}}{\hat{\rho}_{0}} - 4\pi G\frac{a^{2}}{k^{2}})\rho_{1} &= 0, \\ \delta &= \frac{\rho_{1}}{\hat{\rho}} \\ \text{SFDM} & \frac{d^{2}\delta}{dt^{2}} + 2H\frac{d\delta}{dt} + \left(v_{q}^{2}\frac{k^{2}}{a^{2}} - 4\pi G\hat{\rho}_{0}\right)\delta = 0, \quad v_{q}^{2} = \frac{\hbar^{2}k^{2}}{4a^{2}m^{2}} \\ \text{CDM} & \frac{d^{2}\delta}{dt^{2}} + 2H\frac{d\delta}{dt} + \left(v_{s}^{2}\frac{k^{2}}{a^{2}} - 4\pi G\hat{\rho}_{0}\right)\delta = 0, \end{split}$$





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The Cosmology



I. Rodriguez, A. Pérez-Lorenzana, E. de la Cruz, Y. Giraud-Héraud and TM. Bosonic Cosmic Dark Matter. arXiv:1110.2751





Scalar Field Fluctuation = Halo



M.Alcubierre, F. S. Guzmán, T. Matos, D. Núñez, L. A. Ureña and P. Wiederhold. Galactic Collapse of Scalar Field Dark Matter. <u>CQG 19(2002)5017</u>. arXiv:gr-qc /0110102.

Pau Amaro-Seoane, Juan Barranco, Argelia Bernal and Luciano Rezzolla. Constraining scalar fields with stellar kinematics and collisional dark matter. JCAP11 (2010)002

J. Balakrishna, E. Seidel and W. Suen. PRD 58(1998)104004

$$m \sim 1 \text{eV} \qquad \square \gg \qquad \lambda = 1 \times 10^{-6}$$

$$M \sim 0.06\sqrt{\lambda} \, \frac{m_{pl}^3}{m^2} \qquad \qquad T_c = \frac{2m}{\sqrt{\lambda}}$$

 $M \sim 10^{14} M_{\odot}$ $T_c \sim 2000 \, eV$



Axions vs SFDM



Barranco and Bernal, PRD 83,(2011)043525







Natural Cut off



Tonatiuh Matos and Luis A. Ureña. Phys Rev. D63, (2001), 063506. Available at: astro-ph/ 0006024





T. Harko & C. G. Bhömer JCAP 0706:025,(2007) arXiv:0705.4158



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6 8 r/kpc





Galaxy Formation A. Bernal, TM, D. Numez. Rev. Mex. A.A. 44, (2008), 149-160





r[Kps]



r[Kpc]

r[Kps]







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Galaxy Formation

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Galaxy Formation



V. H. Robles and TM. MNRAS 392, (2012) in press. arXiv:1201.3032









Luis A. Medina and Tonatiuh Matos. MNRAS (2014) in press





Scalar Field Fluctuation = Halo



Luis A. Medina and Tonatiuh Matos. MNRAS (2014) in press

















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Hsi-Yu Schive, Tzihong Chiueh and Tom Broadhurst. NATURE PHYSICS, 10 (2014), 246 a ψDM













Hsi-Yu Schive, Tzihong Chiueh and Tom Broadhurst. NATURE PHYSICS, 10 (2014), 246







Hsi-Yu Schive, Tzihong Chiueh and Tom Broadhurst. NATURE PHYSICS, 10 (2014), 246



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Galactic Dynamics



Victor H. Robles, V. Lora, T. Matos & F. J. Sánchez-Salcedo arXiv:1404.3424









Abril Suarez, Victor H. Robles, Tonatiuh Matos . A Review on the Scalar Field/ Bose-Einstein Condensate Dark Matter Model Astrophysics and Space Science Proceedings 38, Chapter 9 (2013) arXiv:1302.0903

Behaves like dust at cosmological lebel

Clusters form by hierarchy

Galaxies form by condensation of the SF

Haloes are BEC drops

MPS has a natural cut off

Same predictions for CMB and MPS

Same predictions for structure formation



Galaxies haloes form erlier and are similar

Galaxies are core

Substructure is restricted







Scalar Field Dark Matter or Ultra light Boson particles are alternative candidates to be the Dark Matter of the Universe







We are in the threshold for a new era in Cosmology, Science and Thought